

INEPT and Signal Enhancement

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The purpose of the INEPT sequence is to transfer the large magnetization from high- γ nuclei, like ^1H or ^{19}F , to low- γ nuclei (labeled 'X'), like ^{13}C and ^{15}N . This transfer significantly improves the signal intensity of 'X' nuclei when measuring their chemical shifts. For example, the magnetic moment and magnetization of ^1H nuclei is roughly 4 times greater than ^{13}C and 10 times greater than ^{15}N . Additionally, the ^1H T_1 is typically much shorter than the T_1 for ^{13}C or ^{15}N , which allows a shorter recycle delay and the collection of a greater number of spectra.

1.1. Theory

The INEPT sequence (Fig. 1.1) transfers magnetization from ^1H nuclei bonded to the 'X' nuclei. The two nuclei must be bonded, or at least connected through intermediary bonds with other atoms, because the nuclei must share a J-coupling.

The INEPT sequence comprises a set of 90° pulses (thin bars), 180° pulses (thick bars) and delays (Δ).

1.1.1. Methine, Amide and the AX Spin Systems

The first 90_x° (^1H) pulse creates $-H_y$ ^1H magnetization. Thereafter, the ^1H 180° pulse in the middle of the two Δ delays acts as a Hahn-echo to refocus the ^1H chemical shifts. Without the accompanying 180° pulse on the 'X' channel, the J-coupling between the ^1H and X nucleus would be refocused as well—and nothing would be accomplished. The 180° pulse on the 'X' channel acts to *cancel* the refocusing effect of the ^1H 180° pulse for the J-coupling.

During the first delta period, the $-H_y$ magnetization evolves under the ^1H chemical shift (ω_H) and the ^1H -X J-coupling. Since the chemical shift and J-coupling Hamiltonians commute, the rotations of each can be conducted sequentially.

In the first step, the operators are propagated for the 90_x° (^1H) pulse and the first ' Δ ' period. We'll use a ^{15}N nucleus as an example of an 'X'

Morris GA, Freeman R. Enhancement of Nuclear Magnetic Resonance Signals by Polarization Transfer. *J Am Chem Soc.* 1979;101(3):760–762.

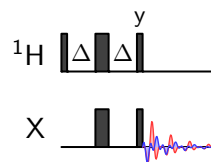


Fig. 1.1. The basic INEPT sequence between nuclei ^1H and X.

A 90_x° pulse rotates the magnetization by 90° with a phase of 'x'.

nucleus.

$$\begin{aligned}
 H_z &\xrightarrow{90^\circ_x(^1H)} -H_y \\
 -H_y &\xrightarrow{\omega_H\Delta} -H_y \cos(\omega_H\Delta) + H_x \sin(\omega_H\Delta) \\
 &\xrightarrow{J_{NH}\Delta} -H_y \cos(\omega_H\Delta) \cos(\pi J_{NH}\Delta) + \\
 &\quad 2H_x N_z \cos(\omega_H\Delta) \sin(\pi J_{NH}\Delta) + \\
 &\quad H_x \sin(\omega_H\Delta) \cos(\pi J_{NH}\Delta) + \\
 &\quad 2H_y N_z \sin(\omega_H\Delta) \sin(\pi J_{NH}\Delta)
 \end{aligned}$$

In the second step, the 180°_x pulses on the 1H and 'X' channels inverts H_y and $2H_x N_z$ terms because these are orthogonal to the x-phase of the pulse.

$$\begin{aligned}
 \xrightarrow{180^\circ_x(^1H), 180^\circ_x(X)} & H_y C_{CS\Delta} C_{J_{NH}\Delta} - 2H_x N_z C_{CS\Delta} S_{J_{NH}\Delta} \\
 & + H_x C_{CS\Delta} C_{J_{NH}\Delta} + 2H_y N_z S_{CS\Delta} S_{J_{NH}\Delta}
 \end{aligned}$$

In the third step, we'll propagate the ' Δ ' delay for the chemical shift first.

$$\begin{aligned}
 \xrightarrow{\omega_H\Delta} & H_y C_{CS\Delta} C_{J_{NH}\Delta} C_{CS\Delta} - H_x C_{CS\Delta} C_{J_{NH}\Delta} S_{CS\Delta} \\
 & - 2H_x N_z C_{CS\Delta} S_{J_{NH}\Delta} C_{CS\Delta} - 2H_y N_z C_{CS\Delta} S_{J_{NH}\Delta} S_{CS\Delta} \\
 & H_x S_{CS\Delta} C_{J_{NH}\Delta} C_{CS\Delta} + H_y S_{CS\Delta} C_{J_{NH}\Delta} S_{CS\Delta} + \\
 & 2H_y N_z S_{CS\Delta} S_{J_{NH}\Delta} C_{CS\Delta} + -2H_x N_z S_{CS\Delta} S_{J_{NH}\Delta} S_{CS\Delta}
 \end{aligned}$$

These terms can be grouped.

$$\begin{aligned}
 = & H_x (S_{CS\Delta} C_{J_{NH}\Delta} C_{CS\Delta} - C_{CS\Delta} C_{J_{NH}\Delta} S_{CS\Delta}) \\
 & + H_y (C_{CS\Delta} C_{J_{NH}\Delta} C_{CS\Delta} + S_{CS\Delta} C_{J_{NH}\Delta} S_{CS\Delta}) \\
 & - 2H_x N_z (C_{CS\Delta} S_{J_{NH}\Delta} C_{CS\Delta} + S_{CS\Delta} S_{J_{NH}\Delta} S_{CS\Delta}) \\
 & + 2H_y N_z (S_{CS\Delta} S_{J_{NH}\Delta} C_{CS\Delta} - C_{CS\Delta} S_{J_{NH}\Delta} S_{CS\Delta})
 \end{aligned}$$

Thereafter, we propagate the ' Δ ' delay for the J_{NH} -coupling and simplify the expression.

$$\begin{aligned}
 \xrightarrow{J_{NH}\Delta} & H_x (S_{CS\Delta} C_{J_{NH}\Delta} C_{CS\Delta} - C_{CS\Delta} C_{J_{NH}\Delta} S_{CS\Delta}) C_{J_{NH}\Delta} \\
 & + 2H_y N_z (S_{CS\Delta} C_{J_{NH}\Delta} C_{CS\Delta} - C_{CS\Delta} C_{J_{NH}\Delta} S_{CS\Delta}) S_{J_{NH}\Delta} \\
 & + H_y (C_{CS\Delta} C_{J_{NH}\Delta} C_{CS\Delta} + S_{CS\Delta} C_{J_{NH}\Delta} S_{CS\Delta}) C_{J_{NH}\Delta} \\
 & - 2H_x N_z (C_{CS\Delta} C_{J_{NH}\Delta} C_{CS\Delta} + S_{CS\Delta} C_{J_{NH}\Delta} S_{CS\Delta}) S_{J_{NH}\Delta} \\
 & - 2H_x N_z (C_{CS\Delta} S_{J_{NH}\Delta} C_{CS\Delta} + S_{CS\Delta} S_{J_{NH}\Delta} S_{CS\Delta}) C_{J_{NH}\Delta} \\
 & - H_y (C_{CS\Delta} S_{J_{NH}\Delta} C_{CS\Delta} + S_{CS\Delta} S_{J_{NH}\Delta} S_{CS\Delta}) S_{J_{NH}\Delta} \\
 & + 2H_y N_z (S_{CS\Delta} S_{J_{NH}\Delta} C_{CS\Delta} - C_{CS\Delta} S_{J_{NH}\Delta} S_{CS\Delta}) C_{J_{NH}\Delta} \\
 & - H_x (S_{CS\Delta} S_{J_{NH}\Delta} C_{CS\Delta} - C_{CS\Delta} S_{J_{NH}\Delta} S_{CS\Delta}) S_{J_{NH}\Delta}
 \end{aligned}$$

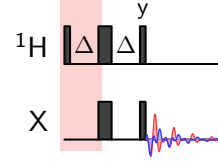


Fig. 1.2. The first step of the INEPT sequence highlighted in red.

The many cosine and sine terms can be abbreviated to simplify their notation in equations.

$$\begin{aligned}
 \cos(\omega_H\Delta) &= C_{CS\Delta} \\
 \cos(\pi J_{NH}\Delta) &= C_{J_{NH}\Delta} \\
 \sin(\omega_H\Delta) &= S_{CS\Delta} \\
 \sin(\pi J_{NH}\Delta) &= S_{J_{NH}\Delta}
 \end{aligned}$$

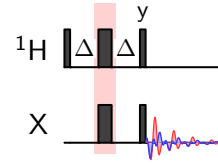


Fig. 1.3. The second step of the INEPT sequence highlighted in red.

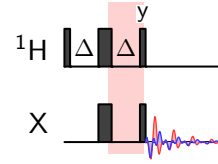


Fig. 1.4. The third step of the INEPT sequence highlighted in red.

$$\begin{aligned}
&= H_x (S_{CS\Delta} C_{J_{NH}\Delta} C_{CS\Delta} C_{J_{NH}\Delta} - C_{CS\Delta} C_{J_{NH}\Delta} S_{CS\Delta} C_{J_{NH}\Delta} \\
&\quad - S_{CS\Delta} S_{J_{NH}\Delta} C_{CS\Delta} S_{J_{NH}\Delta} + C_{CS\Delta} S_{J_{NH}\Delta} S_{CS\Delta} S_{J_{NH}\Delta}) \\
&\quad H_y (C_{CS\Delta} C_{J_{NH}\Delta} C_{CS\Delta} C_{J_{NH}\Delta} + S_{CS\Delta} C_{J_{NH}\Delta} S_{CS\Delta} C_{J_{NH}\Delta} \\
&\quad - C_{CS\Delta} S_{J_{NH}\Delta} C_{CS\Delta} S_{J_{NH}\Delta} - S_{CS\Delta} S_{J_{NH}\Delta} S_{CS\Delta} S_{J_{NH}\Delta}) \\
&\quad - 2H_x N_z (C_{CS\Delta} C_{J_{NH}\Delta} C_{CS\Delta} S_{J_{NH}\Delta} + S_{CS\Delta} C_{J_{NH}\Delta} S_{CS\Delta} S_{J_{NH}\Delta} \\
&\quad + C_{CS\Delta} S_{J_{NH}\Delta} C_{CS\Delta} C_{J_{NH}\Delta} + S_{CS\Delta} S_{J_{NH}\Delta} S_{CS\Delta} C_{J_{NH}\Delta}) \\
&\quad 2H_y N_z (S_{CS\Delta} C_{J_{NH}\Delta} C_{CS\Delta} S_{J_{NH}\Delta} - C_{CS\Delta} C_{J_{NH}\Delta} S_{CS\Delta} S_{J_{NH}\Delta} \\
&\quad + S_{CS\Delta} S_{J_{NH}\Delta} C_{CS\Delta} C_{J_{NH}\Delta} - C_{CS\Delta} S_{J_{NH}\Delta} S_{CS\Delta} C_{J_{NH}\Delta}) \\
&= -2H_x N_z \sin(2\pi J_{NH}\Delta) + H_y \cos(2\pi J_{NH}\Delta)
\end{aligned}$$